Solving three-dimensional Volterra integral equation by the reduced differential transform method

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Abstract

In this article, the results of two-dimensional reduced differential transform method is extended to three-dimensional case for solving three dimensional Volterra integral equation. Using the described method, the exact solution can be obtained after a few number of iterations. Moreover, examples on both linear and nonlinear Volterra integral equation are carried out to illustrate the efficiency and the accuracy of the presented method.

Keywords: Volterra integral equation, Differential transform, Reduce differential transform.

1. Introduction

Many problems in many branches of science such as engineering, physics and other disciplines can be modeled as a Volterra integral equation of the second kind. In literature, different methods have been used to solve one and two-dimensional Volterra integral equation see for example [2, 7, 10, 11, 14] and the references in [3, 13]. For three-dimensional integral equations, in [4] the author used the three-dimensional differential transform method and authors in [8] applied the block-pulse functions methods on three-dimensional nonlinear mixed Volterra-Fredholm integral equation. Recently, the differential transform method is modified to the so-called Reduced Differential Transform Method to solve Volterra integral equation [1].

The aim of the presented article is to solve three-dimensional Volterra integral equation using the reduced differential transform method. So, we consider the 3-dimensional Volterra integral equation of the form:

\[ u(x, y, t) = f(x, y, t) + \int_0^t \int_0^x \int_0^y K(x, y, t, \omega, \nu, \tau)u(\omega, \nu, \tau)d\omega d\nu d\tau, \tag{1} \]

where \( u(x, y, t) \) is the unknown function, \( m \) is a positive integer and the functions \( K \) and \( f \) are analytic in the domain of interest.

2. Reduced differential transform method

In this section, we present basic definitions and operations of the reduced differential transform method, for more details see [12] and the references therein.

Now, assume that the function of three variables \( w(x, y, t) \) can be written as a multiple of two functions as follows:

\[ w(x, y, t) = \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} F(i, j)x^i y^j G(k)t^k = \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} W(i, j)x^i y^j t^k, \tag{2} \]

where the function \( W(i, j) = F(i, j)G(k) \) is called the spectrum of \( w(x, y, t) \).

Definition 2.1. Let \( w(x, y, t) \) be an analytic function in the domain of interest, the reduced differential transform function is

\[ W_k(x, y) = \frac{1}{k!} \left[ \frac{\partial^k}{\partial t^k} w(x, y, t) \right]_{t=0} \quad . \tag{3} \]

Definition 2.2. The differential inverse reduced transform of \( W_k(x, y) \) is defined by

\[ w(x, y, t) = \sum_{k=0}^{\infty} W_k(x, y)(t-t_0)^k = \sum_{k=0}^{\infty} \frac{1}{k!} \left[ \frac{\partial^k}{\partial t^k} w(x, y, t) \right]_{t=t_0} (t-t_0)^k . \tag{4} \]

In fact, the function \( w(x, y, t) \) can be written in a finite series as follows

\[ w_n(x, y, t) = \sum_{k=0}^{n} W_k(x, y)(t-t_0)^k + R_n(x, y, t), \tag{5} \]

the tail function \( R_n(x, y, t) \) is negligibly small.

Table 1: Basic operations of the reduced differential transfom method

<table>
<thead>
<tr>
<th>Original function</th>
<th>Reduced differential transformed</th>
</tr>
</thead>
<tbody>
<tr>
<td>( w(x, y, t) = u(x, y, t)w(x, y, t) )</td>
<td>( W_k(x, y, t) = \sum_{j=0}^{\infty} U_j(x, y)W_{k-j}(x, y) )</td>
</tr>
<tr>
<td>( w(x, y, t) = u(x, y, t) \pm v(x, y, t) )</td>
<td>( W_k(x, y, t) = \sum_{j=0}^{\infty} U_j(x, y) \pm V_{k-j}(x, y) )</td>
</tr>
<tr>
<td>( \frac{\partial^m}{\partial x^m} w(x, y, t) )</td>
<td>( \frac{\partial^m}{\partial x^m} W_k(x, y, t) )</td>
</tr>
<tr>
<td>( \sin(\alpha x + \beta y + \omega t) )</td>
<td>( \frac{\partial^m}{\partial x^m} \sin(\frac{\alpha t}{\omega} + \beta t) )</td>
</tr>
<tr>
<td>( \cos(\alpha x + \beta y + \omega t) )</td>
<td>( \frac{\partial^m}{\partial x^m} \cos(\frac{\alpha t}{\omega} + \beta t) )</td>
</tr>
<tr>
<td>( e^{\alpha x + \beta y + \omega t} )</td>
<td>( e^{\alpha x + \beta y + \omega t} )</td>
</tr>
<tr>
<td>( x^q y^q t^k )</td>
<td>( x^q y^q t^k )</td>
</tr>
</tbody>
</table>

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3. Main results

Suppose that the functions \(W_k(x,y), \ U_k(x,y), \ G_k(x,y)\) and \(V_k(x,y)\) are the reduced differential transform functions of \(w(x,y,t), u(x,y,t), \ g(x,y,t)\) and \(v(x,y,t)\) respectively.

Theorem 3.1. If \(w(x,y,t) = \int_{0}^{x} \int_{0}^{y} u(z, \omega, \tau) dz d\omega d\tau\) then
\[
W_k(x,y) = \frac{1}{k} \int_{0}^{x} \int_{0}^{y} U_{k-1}(\omega, \tau) d\omega d\tau.
\]

Proof. The \(k^{th}\) partial derivative of the function \(w(x,y,t)\) is
\[
\frac{\partial^k}{\partial t^k} w(x,y,t) = \int_{0}^{x} \int_{0}^{y} u(z, \omega, \tau) dz d\omega.
\]
The result can be easily deduced from equation (3).

Theorem 3.2. Let \(w(x,y,t) = \int_{0}^{x} \int_{0}^{y} u(z, \omega, \tau) v(x,y) dz d\omega d\tau\) then
The reduced differential transform function of \(w(x,y,t)\) is
\[
W_k(x,y) = \frac{1}{k} \int_{0}^{x} \int_{0}^{y} \sum_{r=0}^{k-1} U_{k-r-1}(\omega, \tau) V_{k-r}(x,y) d\omega d\tau.
\]

Proof. From Leibnitz formula the \(k^{th}\) partial derivative of \(w(x,y,t)\) is
\[
\int_{0}^{x} \int_{0}^{y} \sum_{r=0}^{k-1} \binom{k-1}{r} U_{k-r-1}(\omega, \tau) V_r(x,y) d\omega d\tau,
\]

table (1) and equation (3) yield the following equation
\[
\frac{\partial^k}{\partial t^k} w(x,y,t) = \int_{0}^{x} \int_{0}^{y} \sum_{r=0}^{k-1} \binom{k-1}{r} U_{k-r-1}(\omega, \tau) V_r(x,y) d\omega d\tau.
\]

Table 3.3. Let \(w(x,y,t) = h(x,y,t) \int_{0}^{x} \int_{0}^{y} u(z, \omega, \tau) dz d\omega d\tau\), then
the reduced differential transform of \(w(x,y,t)\) is
\[
W_k(x,y) = \frac{1}{k} \int_{0}^{x} \int_{0}^{y} H_r(x,y) U_{k-r-1}(\omega, \tau) d\omega d\tau.
\]

Proof. The \(k^{th}\) partial derivative of \(w(x,y,t)\) with respect to \(t\) is
\[
\frac{\partial^k}{\partial t^k} w(x,y,t) = \sum_{r=0}^{k} \binom{k}{r} \frac{\partial^r}{\partial t^r} h(x,y,t) \int_{0}^{x} \int_{0}^{y} \frac{\partial^{k-r-1}}{\partial t^{k-r-1}} u(\omega, \tau, t) d\omega d\tau.
\]

On the other hand, \(\frac{\partial^{k-r-1}}{\partial t^{k-r-1}} u(\omega, \tau, t) \int_{0}^{x} \int_{0}^{y} u(z, \omega, \tau) dz d\omega d\tau\) = 0 for \(k = r\).
Thus, from equation (3) we have \(\frac{\partial^k}{\partial t^k} w(x,y,t) = k! W_k(x,y)\) and hence,
\[
\frac{\partial^k}{\partial t^k} w(x,y,t) = \sum_{r=0}^{k-1} \binom{k}{r} \frac{\partial^r}{\partial t^r} h(x,y,t) \int_{0}^{x} \int_{0}^{y} U_{k-r-1}(\omega, \tau) d\omega d\tau.
\]

Theorem 3.4. If \(w(x,y,t) = \int_{0}^{x} \int_{0}^{y} u(z, \omega, \tau) dz d\omega d\tau\), then
\[
U_k(x,y) = \sum_{r=0}^{k} \binom{k}{r} \frac{\partial^{r+1}}{\partial t^{r+1}} w(x,y,t) \frac{\partial^k}{\partial t^k} u(\omega, \tau, t).
\]

Proof. First write the function \(u(x,y,t) = \frac{\partial^k}{\partial t^k} w(x,y,t)\) then we differentiate partially \(k\) times with respect to \(t\) to get the following equation
\[
\frac{\partial^k}{\partial t^k} u(x,y,t) = \sum_{r=0}^{k} \binom{k}{r} \frac{\partial^{r+1}}{\partial t^{r+1}} w(x,y,t) \frac{\partial^k}{\partial t^k} u(\omega, \tau, t).
\]
Therefore, from equation (3) we have
\[
k! U_k(x,y) = \sum_{r=0}^{k} \binom{k}{r} \frac{\partial^{r+1}}{\partial t^{r+1}} w(x,y,t) \frac{\partial^k}{\partial t^k} u(\omega, \tau, t).
\]

4. Numerical examples

In this section, we apply RDTM method on several examples of linear and nonlinear Volterra integral equation. Then we close this section by an example which is solved in [4] by DTM method and present the table of absolute error at some particular points to compare between the two methods.

Example 4.1. Consider the integral equation
\[
\int_{0}^{x} \int_{0}^{y} u(z, \omega, \tau) dz d\omega d\tau = \int_{0}^{x} \int_{0}^{y} \int_{0}^{z} 2x \sin t d\tau d\omega.
\]
The exact solution is \(u(x,y,t) = x \cos t\). First, let \(f(x,y) = x \cos t\) and \(v(x,y,t) = x^2 t^2\).

Elementary calculations lead to \(F_0(x,y) = x, F_1(x,y) = -\frac{x^2}{2}, \ F_2(x,y) = \frac{x^3}{6}, \ F_3(x,y) = -\frac{x^4}{24}, \ldots\) and \(V_0(x,y) = x^2, \ V_1(x,y) = 0\) for \(k \geq 1\). Also, \(U_0(x,y) = F_0(x,y) = x\). Now, from theorem (3.2) \(U_k(x,y) = F_k(x,y) + \frac{1}{k!} \sum_{r=0}^{k-1} U_{k-r-1}(\omega, \tau) V_r(x,y) d\omega d\tau, k \geq 1\).

So, exact calculation on the last integral equation will produce the following formula:
\[
U_k(x,y) = \left\{ \begin{array}{ll}
\frac{(-1)^{k+1}}{k}, & \text{k is even} \\
0, & \text{k is odd}
\end{array} \right.
\]

Therefore,
\[
u(x,y,t) = \sum_{k=0}^{\infty} U_k(x,y) t^k = x \sum_{k=0}^{\infty} \left\{ \frac{(-1)^{k+1}}{(2k)!} \right\} t^k = x \cos t.
\]

Example 4.2. Consider the nonlinear integral equation
\[
u(x,y,t) = \int_{0}^{x} \int_{0}^{y} u^3(z, \omega, \tau) dz d\omega d\tau.
\]
The exact solution to this integral equation is \(u(x,y,t) = x y t\).
It is clear that \( U_0(x,y) = F_0(x,y) = 0, k \neq 1 \) and \( F_1(x,y) = xy \). Also, from theorem(3.2) we have

\[
U_k(x,y) = F_k(x,y) + \frac{1}{k}\int_0^1 \sum_{r=0}^{k-1} U_r(\omega, \tau) U_{k-r-1}(\omega, \tau) d\omega d\tau, k \geq 1.
\]

Hence \( U_1(x,y) = xy \) and \( U_k(x,y) = 0 \) for \( k \neq 1 \).

**Example 4.3.** Consider the following integral equation

\[
w(x,y,t) = g(x,y,t) + w(x,y,t) \text{ where } v(x,y,t) = 1 + \cos x.
\]

It is obvious that \( U_0(x,y) = G_0(x,y) = 2(x+y) \). Also, from table (1), \( G_1(x,y) = -\frac{xy + y^2}{2} \), and for \( k \geq 2 \)

\[
G_k(x,y) = \begin{cases} \frac{(-1)^{k+1}(x+y)}{k}, & \text{even} \\ 0, & \text{odd} \end{cases}
\]

On the other hand, \( V_0(x,y) = 2 \) and for \( k \geq 1 \),

\[
V_k(x,y) = \begin{cases} \frac{\pi}{2} \cos \frac{k \pi}{2}, & \text{even} \\ 0, & \text{odd} \end{cases}
\]

Using (9), \( W_k(x,y) = \frac{x^k y^{k+1}}{k!} \), \( W_2(x,y) = 0 \). But since \( U_k(x,y) = G_k(x,y) + W_k(x,y) \) we have for \( k \geq 1 \),

\[
U_k(x,y) = \begin{cases} \frac{(-1)^{k+1}(x+y)}{k}, & \text{even} \\ 0, & \text{odd} \end{cases}
\]

Therefore,

\[
u(x,y,t) = 2(x+y) + \frac{-x^2}{2} + \frac{y^2}{4} + \frac{(x+y)^4}{4!} + \cdots
\]

\[
=(x+y)(1 + \frac{1}{2} + \frac{1}{4} + \frac{1}{4!} + \cdots)
\]

\[
=(x+y)(1 + \cos t).
\]

**Example 4.4.** Consider the following integral equation

\[
u(x,y,t) = g(x,y,t) + \int_0^y \int_0^x u(\omega,\tau) d\omega d\tau, 0 \leq x \leq y \leq 1
\]

and \( g(x,y,t) = e^{\epsilon+y-1} \). The exact solution is \( e^{\epsilon+y} \).

It is clear that \( U_0(x,y) = G_0(x,y) = e^{\epsilon+y} \). Also, from table (1), \( G_2(x,y) = \frac{1}{2}(e^\epsilon + e^y - 1) \). Theorem (3.1) implies

\[
U_k(x,y) = G_k(x,y) + \frac{1}{k} \int_0^y \int_0^x U_{k-1}(\omega, \tau) d\omega d\tau, k = 1, 2, \ldots
\]

Using equation (11) recursively one can obtain the general formula of \( U_k(x,y) \) = \( e^{\epsilon+y} \) and hence \( u(x,y,t) = e^{\epsilon+y} \).

The exact solution and absolute error of the iterative solutions \( u_2(x,y,t) \) and \( u_3(x,y,t) \) obtained by RDTM at some test points \((x,y,t)\) are calculated in tables(2) and (3). Comparing the results obtained by RDTM with those obtained by DTM method (see example 3.3 [4]), we conclude that RDTM minimizes the number of iterations to reach the exact solution. Also, the approximated solution approaches rapidly to the exact solution.

<table>
<thead>
<tr>
<th>Table 2: The Exact Solution</th>
</tr>
</thead>
<tbody>
<tr>
<td>( x )</td>
</tr>
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<td>0.1</td>
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<tr>
<td>0.01</td>
</tr>
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<tr>
<td>0.00001</td>
</tr>
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<td>0.000001</td>
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</tbody>
</table>

<table>
<thead>
<tr>
<th>Table 3: The Absolute Error</th>
</tr>
</thead>
<tbody>
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<td>( x )</td>
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<tr>
<td>0.00000001</td>
</tr>
</tbody>
</table>

5. Conclusion

In this paper, the three dimensional linear and nonlinear integral equations are solved by using RDTM. It is worth noting that RDTM does not require complex computational work like DTM. It can be easily implemented, its convergence is rapid and its approximation is accurate. In general, it can be concluded that RDTM is a powerful tool for solving many linear and nonlinear three dimensional integral equations.

Acknowledgement

The authors would like to thank Mr. Hayel Al-Shraydeh for his support in improving the print of the manuscript.

References


