# Blow-up result in a Cauchy problem for the nonlinear viscoelastic Petrovsky equation 

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#### Abstract

In this paper, we consider a Cauchy problem for the nonlinear viscoelastic Petrovsky equation. We obtain the blow up of solutions by applying a lemma due to Zhou.


Keywords: Blow Up; Cauchy Problem; Nonlinear Viscoelastic Petrovsky Equation.

## 1. Introduction

In [5], Li et al. considered the following nonlinear viscoelastic Petrovsky problem
$\left\{\begin{array}{lr}u_{t t}+\Delta^{2} u-\int_{0}^{t} g(t-\tau) \Delta^{2} u(t, \tau) d \tau-\Delta u_{t t}-\Delta u_{t}+\left|u_{t}\right|^{m-1} u_{t}=|u|^{p-1} u, x \in \Omega, t>0, \\ u(x, t)=\frac{\partial u(x, t)}{\partial v}=0, & x \in \Omega, \\ u(x, 0)=u_{0}(x), u_{t}(x, 0)=u_{1}(x), & x \in \Omega,\end{array}\right.$
where $\Omega$ is a bounded domain in $R^{n}(n \geq 1)$ with a smooth boundary $\partial \Omega, m, p \geq 1 ; v$ is the unit outer normal on $\partial \Omega$; and g is a nonnegative memory term. They established some asymptotic behavior and blow up results for solutions with positive initial energy.
Guesmia [3] studied the problem
$u_{t t}+\Delta^{2} u+q(x) u+g\left(u_{t}\right)=0$,
where $q: \Omega \rightarrow R^{+}$is a bounded function? Under some assumptions, he showed the solution of (2) decay results by using the semigroup method. In [7], Messaoudi investigated the semilinear Petrovsky equation
$u_{t t}+\Delta^{2} u+\left|u_{t}\right|^{m-1} u_{t}=|u|^{p-1} u$.
He showed that the solution blows up in finite time if $p>m$ and while it exists globally if $p \leq m$. in [9], Wu and Tsai showed that the solution of (3) is global under some conditions. Also, Chen and Zhou [2] studied the blow up of the solution of (3).
Recently, Li et al. [6] considered the following Petrovsky equation
$u_{t t}+\Delta^{2} u-\Delta u_{t}+\left|u_{t}\right|^{m-1} u_{t}=|u|^{p-1} u$.
The authors obtained global existence, decay and blow up of the solution. Very recently, Pişkin and Polat [8] studied the decay of the solution of the problem (4).
In this paper, our aim is to extend the result of [5], established in bounded domains, to the problem in unbounded domains. Namely, we consider the following Cauchy problem

$$
\left\{\begin{array}{lc}
u_{t t}+\Delta^{2} u-\int_{0}^{t} g(t-\tau) \Delta^{2} u(t, \tau) d \tau-\Delta u_{t t}-\Delta u_{t}+\left|u_{t}\right|^{m-1} u_{t}=|u|^{p-1} u, x \in R^{n}, t>0  \tag{5}\\
u(x, 0)=u_{0}(x), u_{t}(x, 0)=u_{1}(x), & x \in R^{n}
\end{array}\right.
$$

where $g, u_{0}, u_{1}$ are functions to be specified later.
This paper is organized as follows. In section 2 , we present some notations, lemmas, and the local existence theorem. In section 3, under suitable conditions on the initial data, we prove a finite time blow up result.

## 2. Preliminary notes

In this section, we give some assumptions and lemmas which will be used throughout this work. Hereafter we denote by $\|\cdot\|$ and $\|\cdot\|_{p}$ the norm of $L^{2}\left(R^{n}\right)$ and $L^{p}\left(R^{n}\right)$, respectively. First, we make the following assumptions
(G) $g: R^{+} \rightarrow R^{+}$is a nonincreasing differentiable function such that
$1-\int_{0}^{\infty} g(\tau) d \tau=l>0, g^{\prime}(t) \leq 0, t \geq 0$.
Next, we state the local existence theorem of the problem (5), which can be established by combining the arguments of [1], [7].

Theorem 1: (Local existence). Suppose that $(G)$ holds, and $1<p<\infty$ if $n=1,2,3,4$, and $1<p<\frac{n}{n-4}$ if $n \geq 5$. Then for any initial data $\left(u_{0}, u_{1}\right) \in H^{2}\left(R^{n}\right) \times H^{1}\left(R^{n}\right)$, with compact support, the problem (5) has a unique local solution
$u \in C\left([0, T) ; H^{2}\left(R^{n}\right)\right) \cap C^{1}\left([0, T) ; L^{2}\left(R^{n}\right)\right)$,
$u_{t} \in L^{2}\left([0, T) ; H^{1}\left(R^{n}\right)\right) \cap L^{2}(\Omega \times[0, T))$
for $T$ small enough.
To obtain the result of this paper, we will introduce the modified energy functional

$$
\begin{equation*}
E(t)=\frac{1}{2}\left\|u_{t}\right\|^{2}+\frac{1}{2}\|\nabla u\|^{2}+\frac{1}{2}\left(1-\int_{0}^{t} g(\tau) d \tau\right)\|\Delta u\|^{2}+\frac{1}{2}(g \circ \Delta u)(t)-\frac{1}{p+1}\|u\|_{p+1}^{p+1}, \tag{6}
\end{equation*}
$$

where
$(g \circ v)(t)=\int_{0}^{t} g(t-\tau)\|v(t)-v(\tau)\| d \tau$.
The next lemma shows that our energy functional (6) is a nonincreasing function along the solution of (5).
Lemma 2: $E(t)$ is a nonincreasing function for $t \geq 0$ and
$E^{\prime}(t)=-\left(\left\|\nabla u_{t}\right\|^{2}+\left\|u_{t}\right\|^{2}\right) \leq 0$.
Proof: By multiplying the equation in (5) by $u_{t}$ and integrating over $R^{n}$, we obtain (7).

## 3. Blow up of solutions

In this section, we shall show that the solution of the problem (5) blow up in finite time, by the similar arguments as in [4]. For the purpose, we give the lemma.

Lemma 3: [10] Suppose that $\psi(t)$ is a twice continuously differentiable function satisfying
$\left\{\begin{array}{l}\psi^{\prime \prime}(t)+\psi^{\prime}(t) \geq C_{0}(t+L)^{\beta} \psi^{1+\alpha}(t), t>0, \\ \psi(0)>0, \psi^{\prime}(0) \geq 0,\end{array}\right.$
where $C_{0}, L>0,-1<\beta \leq 0, \alpha>0$ are constants. Then, $\psi(t)$ blow up in finite time.

Theorem 4: Suppose that ( $G$ ) holds, and $1<p<\infty$ if $n=1,2,3,4$, and $1<p<\frac{n}{n-4}$ if $n \geq 5$. Assume further that $\int_{0}^{t} g(\tau) d \tau<\frac{p^{2}-1}{p^{2}}$.

Then for any initial data $\left(u_{0}, u_{1}\right) \in H^{2}\left(R^{n}\right) \times H^{1}\left(R^{n}\right)$, with compact support, satisfying
$E(0) \leq 0, \int_{R} n u_{0} u_{1} d x \geq 0,\left\|\nabla u_{0}\right\|^{2} \leq\left\|u_{0}\right\|^{2}$,
Then the corresponding solution blows up in finite time. In other words, there exists a positive constant $T^{*}$ such that $\lim _{t \rightarrow T^{*}}\|u\|^{2}=\infty$.
Proof: By multiplying the equation in (5) by $u_{t}$ and integrating over $R^{n}$, using integrating by parts, we obtain
$\frac{1}{2} \frac{d}{d t}\left(\int_{R^{n}}\left|u_{t}\right|^{2} d x+\int_{R} n|\Delta u|^{2} d x+\int_{R} n\left|\nabla u_{t}\right|^{2} d x\right)+\int_{R} n\left|\nabla u_{t}\right|^{2} d x+\int_{R} n\left|u_{t}\right|^{2} d x-\int_{0}^{t} g(t-\tau) \int_{R} n \Delta u(\tau) \Delta u_{t}(t) d x d \tau$ $=\frac{1}{p+1} \frac{d}{d t} \int_{R^{n}}|u|^{p+1} d x$,
the last term on the left side of (9) can be estimated as follows

$$
\begin{align*}
\int_{0}^{t} g(t-\tau) \int_{R}^{n \Delta u \Delta u_{t} d x d \tau=} & \int_{0}^{t} g(t-\tau) \int_{R} n[\Delta u(\tau)-\Delta u(t)] \Delta u_{t}(t) d x d \tau+\int_{0}^{t} g(t-\tau) \int_{R}^{n \Delta u(t) \Delta u_{t}(t) d x d \tau} \\
= & -\frac{1}{2} \int_{0}^{t} g(t-\tau) \frac{d}{d t} \int_{R}^{n}[\Delta u(\tau)-\Delta u(t)]^{2} d x d \tau+\int_{0}^{t} g(t-\tau)\left(\frac{1}{2} \frac{d}{d t} \int_{R} n|\Delta u(t)|^{2} d x\right) d \tau \\
= & -\frac{1}{2} \frac{d}{d t}\left[\int_{0}^{t} g(t-\tau) \int_{R} n[\Delta u(\tau)-\Delta u(t)]^{2} d x d \tau\right]+\frac{1}{2} \int_{0}^{t} g^{\prime}(t-\tau)\left(\int_{R} n[\Delta u(\tau)-\Delta u(t)]^{2} d x\right) d \tau \\
& +\frac{1}{2} \frac{d}{d t}\left[\int_{0}^{t} g(\tau) \int_{R} n|\Delta u(t)|^{2} d x d \tau\right]-\frac{1}{2} g(t) \int_{R} n|\Delta u(t)|^{2} d x . \tag{10}
\end{align*}
$$

Inserting (10) into (9), to get
$\frac{1}{2} \frac{d}{d t}\left[\int_{R} n\left|u_{t}\right|^{2} d x+\int_{R} n|\Delta u|^{2} d x+\int_{R} n\left|\nabla u_{t}\right|^{2} d x-\frac{2}{p+1} \int_{R} n|u|^{p+1} d x\right.$
$\left.+\int_{0}^{t} g(t-\tau) \int_{R} n[\Delta u(\tau)-\Delta u(t)]^{2} d x d \tau-\int_{0}^{t} g(\tau) \int_{R} n|\Delta u(t)|^{2} d x d \tau\right]$
$=\frac{1}{2} \int_{0}^{t} g^{\prime}(t-\tau) \int_{R} n[\Delta u(\tau)-\Delta u(t)]^{2} d x d \tau-\frac{1}{2} g(t) \int_{R} n|\Delta u(t)|^{2} d x$
$-\int_{R} n\left|u_{t}\right|^{2} d x-\int_{R} n\left|\nabla u_{t}\right|^{2} d x$.
To apply Lemma 3, we define
$\psi(t)=\frac{1}{2} \int_{R} n\left(|u|^{2}+|\nabla u|^{2}\right) d x$.
Therefore

$$
\begin{equation*}
\psi^{\prime}(t)=\int_{R} n\left(u u_{t}+\nabla u \nabla u_{t}\right) d x, \tag{12}
\end{equation*}
$$

and

$$
\begin{equation*}
\psi^{\prime \prime}(t)=\int_{R} n\left(u u_{t t}+\left|u_{t}\right|^{2}+\nabla u \nabla u_{t t}+\left|\nabla u_{t}\right|^{2}\right) d x \tag{13}
\end{equation*}
$$

Then, eq (5) is used to estimate (13) as follows

$$
\begin{align*}
\psi^{\prime \prime}(t)= & \int_{R} n\left(u u_{t t}+\left|u_{t}\right|^{2}+\nabla u \nabla u_{t t}+\left|\nabla u_{t}\right|^{2}\right) d x \\
= & -\int_{R} n|\Delta u|^{2} d x-\int_{0}^{t} \Delta u(t) \int_{R} n g(t-\tau) \Delta u(\tau) d x d \tau-\int_{R} n \nabla u \nabla u_{t} d x \\
& -\int_{R} n u u_{t} d x+\int_{R} n u|u|^{p-1} u d x+\int_{R} n\left(\left|u_{t}\right|^{2}+\left|\nabla u_{t}\right|^{2}\right) d x . \tag{14}
\end{align*}
$$

On using
$\int_{0}^{t} \Delta u(t) \int_{R} n g(t-\tau) \Delta u(\tau) d x d \tau=\int_{0}^{t} g(t-\tau) \int_{R} n \Delta u(t)[\Delta u(\tau)-\Delta u(t)] d x d \tau+\int_{0}^{t} g(\tau) d \tau \int_{R} n|\Delta u(t)|^{2} d x$.
Eq. (14) becomes

$$
\begin{align*}
\psi^{"}(t)= & -\left(1-\int_{0}^{t} g(\tau) d \tau\right) \int_{R} n|\Delta u(t)|^{2} d x-\int_{0}^{t} g(t-\tau) \int_{R} n \Delta u(t)[\Delta u(\tau)-\Delta u(t)] d x d \tau \\
& -\int_{R} n \nabla u \nabla u_{t} d x-\int_{R} n u u_{t} d x+\int_{R} n|u|^{p+1} d x+\int_{R} n\left(\left|u_{t}\right|^{2}+\left|\nabla u_{t}\right|^{2}\right) d x . \tag{15}
\end{align*}
$$

We then use Young inequality to estimate the second term in (15). Namely,
$\int_{R}^{n \Delta u(t)} \int_{0}^{t} g(t-\tau)[\Delta u(\tau)-\Delta u(t)] d \tau d x \leq \delta \int_{R} n|\Delta u(t)|^{2} d x+\frac{1}{4 \delta} \int_{R} n\left|\int_{0}^{t} g(t-\tau)(\Delta u(\tau)-\Delta u(\tau)) d \tau\right|^{2} d x$

$$
\begin{equation*}
\leq \delta \int_{R} n|\Delta u(t)|^{2} d x+\frac{1}{4 \delta}\left(\int_{0}^{t} g(\tau) d \tau\right)(g \mathrm{o} \Delta u) \tag{16}
\end{equation*}
$$

By combining (15) and (16), we get

$$
\begin{align*}
\psi^{"}(t) \geq & -\left(1-\int_{0}^{t} g(\tau) d \tau\right) \int_{R} n|\Delta u(t)|^{2} d x-\delta \int_{R} n|\Delta u(t)|^{2} d x-\frac{1}{4 \delta}\left(\int_{0}^{t} g(\tau) d \tau\right)(g \text { o } \Delta u) \\
& -\int_{R} n \nabla u \nabla u_{t} d x-\int_{R} n u u_{t} d x+\int_{R} n|u|^{p+1} d x+\int_{R} n\left(\left|u_{t}\right|^{2}+\left|\nabla u_{t}\right|^{2}\right) d x . \tag{17}
\end{align*}
$$

From (12), (13) and (17), we obtain

$$
\begin{align*}
\psi^{\prime \prime}(t)+\psi^{\prime}(t) \geq & -\left(1+\delta-\int_{0}^{t} g(\tau) d \tau\right) \int_{R} n|\Delta u(t)|^{2} d x-\frac{1}{4 \delta}\left(\int_{0}^{t} g(\tau) d \tau\right)(g \circ \Delta u) \\
& +\int_{R} n|u|^{p+1} d x+\int_{R} n\left(\left|u_{t}\right|^{2}+\left|\nabla u_{t}\right|^{2}\right) d x . \tag{18}
\end{align*}
$$

Now, we exploit (6) to substitute for $(g \circ \Delta u)(t)$;
$(g \mathrm{o} \Delta u)(t)=2 E(t)-\left\|u_{t}\right\|^{2}-\|\nabla u\|^{2}-\left(1-\int_{0}^{t} g(\tau) d \tau\right)\|\Delta u\|^{2}+\frac{2}{p+1}\|u\|_{p+1}^{p+1}$.
Thus (18) takes the form

$$
\begin{align*}
\psi^{\prime \prime}(t)+\psi^{\prime}(t) \geq & -\frac{1}{2 \delta}\left(\int_{0}^{t} g(\tau) d \tau\right) E(t)+\left[\frac{1}{4 \delta}\left(\int_{0}^{t} g(\tau) d \tau\right)\left(1-\int_{0}^{t} g(\tau) d \tau\right)-\left(1+\delta-\int_{0}^{t} g(\tau) d \tau\right)\right]\|\Delta u\|^{2} \\
& +\left[1+\frac{1}{4 \delta}\left(\int_{0}^{t} g(\tau) d \tau\right)\right]\left\|u_{t}\right\|^{2}+\frac{1}{4 \delta}\left(\int_{0}^{t} g(\tau) d \tau\right)\|\nabla u\|^{2}+\left\|\nabla u_{t}\right\|^{2}+\left[1-\frac{1}{4 \delta}\left(\int_{0}^{t} g(\tau) d \tau\right) \frac{2}{p+1}\right]\|u\|_{p+1}^{p+1} . \tag{19}
\end{align*}
$$

At this point we choose $\delta>0$ so that
$\frac{1}{4 \delta}\left(\int_{0}^{t} g(\tau) d \tau\right)\left(1-\int_{0}^{t} g(\tau) d \tau\right)-\left(1+\delta-\int_{0}^{t} g(\tau) d \tau\right) \geq 0$
and
$1-\frac{1}{4 \delta}\left(\int_{0}^{t} g(\tau) d \tau\right) \frac{2}{p+1}>0$.
This is, of course, possible by (8). We then conclude, from (19), that
$\psi^{\prime \prime}(t)+\psi^{\prime}(t) \geq \gamma\|u\|_{p+1}^{p+1}$.
Now, we use Hölder inequality to estimate $\|u\|_{p+1}^{p+1}$ as follows
$\int_{R} n|u|^{2} d x \leq\left(\int_{R} n|u|^{p+1} d x\right)^{\frac{2}{p+1}}\left(\int_{B(t+L)^{1 d x}}\right)^{\frac{p-1}{p+1}}$,
where $L>0$ is such that
$\operatorname{supp}\left\{u_{0}(x), u_{1}(x)\right\} \subset B(L)$,
and $B(t+L)$ is the ball, with radius $t+L$, centered at the origin. If we call $W_{n}$ the volume of the unit ball then

$$
\begin{equation*}
\int_{R} n|u|^{p+1} d x \geq\left(\int_{R} n|u|^{2} d x\right)^{\frac{p+1}{2}}\left(W_{n}(t+L)^{n}\right)^{\frac{1-p}{2}} . \tag{21}
\end{equation*}
$$

From the definition of $\psi(t)$, we get

$$
\begin{align*}
{[2 \psi(t)]^{\frac{p+1}{2}} } & =\left[\int_{R} n\left(|u|^{2}+|\nabla u|^{2}\right) d x\right]^{\frac{p+1}{2}} \\
& \leq 2^{\frac{p-1}{2}}\left[\left(\int_{R} n|u|^{2} d x\right)^{\frac{p+1}{2}}+\left(\int_{R^{n}}|\nabla u|^{2} d x\right)^{\frac{p+1}{2}}\right] \tag{22}
\end{align*}
$$

Combining (20)-(21), we have
$\psi^{\prime \prime}(t)+\psi^{\prime}(t) \geq \gamma\left[2 \psi^{\frac{p+1}{2}}(t)-\left(\|\nabla u\|^{2}\right)^{\frac{p+1}{2}}\right]\left(W_{n}(t+L)^{n}\right)^{\frac{1-p}{2}}$.
From assumptions of Theorem, we deduce by continuity that there exists $T * \leq T$ such that
$\|\nabla u\|^{2}-\|u\|^{2} \leq 0, \forall t \in[0, T *)$,

$$
\psi^{\frac{p+1}{2}}(t)-\left(\|\nabla u\|^{2}\right)^{\frac{p+1}{2}} \geq 0 .
$$

Consequently, (22) implies that

$$
\psi^{\prime \prime}(t)+\psi^{\prime}(t) \geq \gamma \psi^{\frac{p+1}{2}}(t)\left(W_{n}\right)^{\frac{1-p}{2}}(t+L)^{n \frac{1-p}{2}} .
$$

It is easy to verify that the requirements of Lemma 3 are satisfied by
$C_{0}=\gamma\left(W_{n}\right)^{\frac{1-p}{2}}>0, \beta=n \frac{1-p}{2}, \alpha=\frac{p+1}{2}>0$.
Therefore $\psi(t)$ blow up in finite.

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